

Finite Element Modelling in Geosciences

FEM in 2D for viscous materials

Introduction to Finite Element Modelling in Geosciences

Summer 2019

Antoine Rozel & Patrick Sanan

Lecture by Marcel Frehner

Recap from 2D continuum mechanics and rheology

Equations

Incompressibility

$$\frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} = 0$$

Conservation of linear momentum (i.e., force balance equation)

$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + F_x = 0$$

$$\frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + F_y = 0$$

Rheology

$$\sigma_{xx} = 4\eta/3 \dot{\epsilon}_{xx} - 2\eta/3 \dot{\epsilon}_{yy} - \tilde{p}$$

$$\sigma_{yy} = -2\eta/3 \dot{\epsilon}_{xx} + 4\eta/3 \dot{\epsilon}_{yy} - \tilde{p}$$

$$\sigma_{xy} = \eta \dot{\gamma}_{xy}$$

Kinematic relation

$$\dot{\epsilon}_{xx} = \partial v_x / \partial x$$

$$\dot{\epsilon}_{yy} = \partial v_y / \partial y$$

$$\dot{\gamma}_{xy} = \partial v_x / \partial y + \partial v_y / \partial x$$

Matrix notation

$$\begin{bmatrix} \frac{\partial}{\partial x} & \frac{\partial}{\partial y} \end{bmatrix} \begin{Bmatrix} v_x \\ v_y \end{Bmatrix} = 0$$

$$\begin{bmatrix} \frac{\partial}{\partial x} & 0 & \frac{\partial}{\partial y} \\ 0 & \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix} \begin{Bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{xy} \end{Bmatrix} + \begin{Bmatrix} F_x \\ F_x \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

$$\begin{Bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{xy} \end{Bmatrix} = \begin{bmatrix} 4\eta/3 & -2\eta/3 & 0 \\ -2\eta/3 & 4\eta/3 & 0 \\ 0 & 0 & \eta \end{bmatrix} \begin{Bmatrix} \dot{\epsilon}_{xx} \\ \dot{\epsilon}_{yy} \\ \dot{\gamma}_{xy} \end{Bmatrix} - \begin{Bmatrix} 1 \\ 1 \\ 0 \end{Bmatrix} \tilde{p}$$

$$\begin{Bmatrix} \dot{\epsilon}_{xx} \\ \dot{\epsilon}_{yy} \\ \dot{\gamma}_{xy} \end{Bmatrix} = \begin{bmatrix} \partial/\partial x & 0 \\ 0 & \partial/\partial y \\ \partial/\partial y & \partial/\partial x \end{bmatrix} \begin{Bmatrix} v_x \\ v_y \end{Bmatrix}$$

Short form

$$\nabla \tilde{\mathbf{v}} = 0$$

$$\tilde{\mathbf{B}}^T \tilde{\boldsymbol{\sigma}} + \tilde{\mathbf{F}} = 0$$

$$\tilde{\boldsymbol{\sigma}} = \mathbf{D} \tilde{\boldsymbol{\epsilon}} - \mathbf{m} \tilde{p}$$

$$\tilde{\boldsymbol{\epsilon}} = \tilde{\mathbf{B}} \tilde{\mathbf{v}}$$

Total system of equations

- Incompressibility

$$\frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} = 0$$

! NEW !

- Short form

$$\tilde{\mathbf{B}}^T \mathbf{D} \tilde{\mathbf{B}} \tilde{\mathbf{v}} - \nabla^T \tilde{p} + \tilde{\mathbf{F}} = \mathbf{0}$$

- Matrix notation

$$\begin{bmatrix} \frac{\partial}{\partial x} & 0 & \frac{\partial}{\partial y} \\ 0 & \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix} \begin{bmatrix} \frac{4}{3}\eta & -\frac{2}{3}\eta & 0 \\ -\frac{2}{3}\eta & \frac{4}{3}\eta & 0 \\ 0 & 0 & \eta \end{bmatrix} \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix} \begin{Bmatrix} v_x \\ v_y \end{Bmatrix} - \begin{bmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{bmatrix} p + \begin{Bmatrix} F_x \\ F_y \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \end{Bmatrix}$$

- Written out (2 equations)

$$\frac{\partial}{\partial x} \left[\frac{4}{3}\eta \frac{\partial v_x}{\partial x} - \frac{2}{3}\eta \frac{\partial v_y}{\partial y} \right] + \frac{\partial}{\partial y} \left[\eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) \right] - \frac{\partial p}{\partial x} + F_x = 0$$

$$\frac{\partial}{\partial y} \left[-\frac{2}{3}\eta \frac{\partial v_x}{\partial x} + \frac{4}{3}\eta \frac{\partial v_y}{\partial y} \right] + \frac{\partial}{\partial x} \left[\eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) \right] - \frac{\partial p}{\partial y} + F_y = 0$$

Force balance equation: Deriving the weak form

- Test functions:

$$\mathbf{N} = [N_1 \quad \dots \quad N_9]$$

- ① Multiplication with test functions:

$$\mathbf{N}^T \frac{\partial}{\partial x} \left[\frac{4}{3} \eta \frac{\partial v_x}{\partial x} - \frac{2}{3} \eta \frac{\partial v_y}{\partial y} \right] + \mathbf{N}^T \frac{\partial}{\partial y} \left[\eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) \right] - \mathbf{N}^T \frac{\partial p}{\partial x} + \mathbf{N}^T F_x = 0$$

$$\mathbf{N}^T \frac{\partial}{\partial y} \left[-\frac{2}{3} \eta \frac{\partial v_x}{\partial x} + \frac{4}{3} \eta \frac{\partial v_y}{\partial y} \right] + \mathbf{N}^T \frac{\partial}{\partial x} \left[\eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) \right] - \mathbf{N}^T \frac{\partial p}{\partial y} + \mathbf{N}^T F_y = 0$$

- ② Spatial integration

$$\int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T \frac{\partial}{\partial x} \left[\frac{4}{3} \eta \frac{\partial v_x}{\partial x} - \frac{2}{3} \eta \frac{\partial v_y}{\partial y} \right] dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T \frac{\partial}{\partial y} \left[\eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) \right] dx dy - \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T \frac{\partial p}{\partial x} dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_x dx dy = 0$$

$$\int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T \frac{\partial}{\partial y} \left[-\frac{2}{3} \eta \frac{\partial v_x}{\partial x} + \frac{4}{3} \eta \frac{\partial v_y}{\partial y} \right] dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T \frac{\partial}{\partial x} \left[\eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) \right] dx dy - \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T \frac{\partial p}{\partial y} dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_y dx dy = 0$$

Force balance equation: Deriving the weak form

③ Integration by parts and dropping arising boundary terms

$$\int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \left[\frac{4}{3} \eta \frac{\partial v_x}{\partial x} - \frac{2}{3} \eta \frac{\partial v_y}{\partial y} \right] dy dx + \int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) dx dy - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} p dx dy = \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_x dx dy$$

$$\int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \left[-\frac{2}{3} \eta \frac{\partial v_x}{\partial x} + \frac{4}{3} \eta \frac{\partial v_y}{\partial y} \right] dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \eta \left(\frac{\partial v_x}{\partial y} + \frac{\partial v_y}{\partial x} \right) dy dx - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial y} p dx dy = \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_y dx dy$$

This is the weak form of the force balance equation!

The finite element approximation

2D elasticity

$$u_x(x, y) \approx N_1 u_{x,1} + N_2 u_{x,2} + N_3 u_{x,3} + N_4 u_{x,4}$$
$$= [N_1 \quad N_2 \quad N_3 \quad N_4] \begin{Bmatrix} u_{x,1} \\ u_{x,2} \\ u_{x,3} \\ u_{x,4} \end{Bmatrix} = \mathbf{N} \mathbf{u}_x$$

$$u_y(x, y) \approx \mathbf{N} \mathbf{u}_y$$

Requirements for the shape functions:

- Equal 1 at one nodal point
- Equal 0 at all other nodal points
- Sum of all shape functions has to be 1 in the whole finite element.

2D viscosity

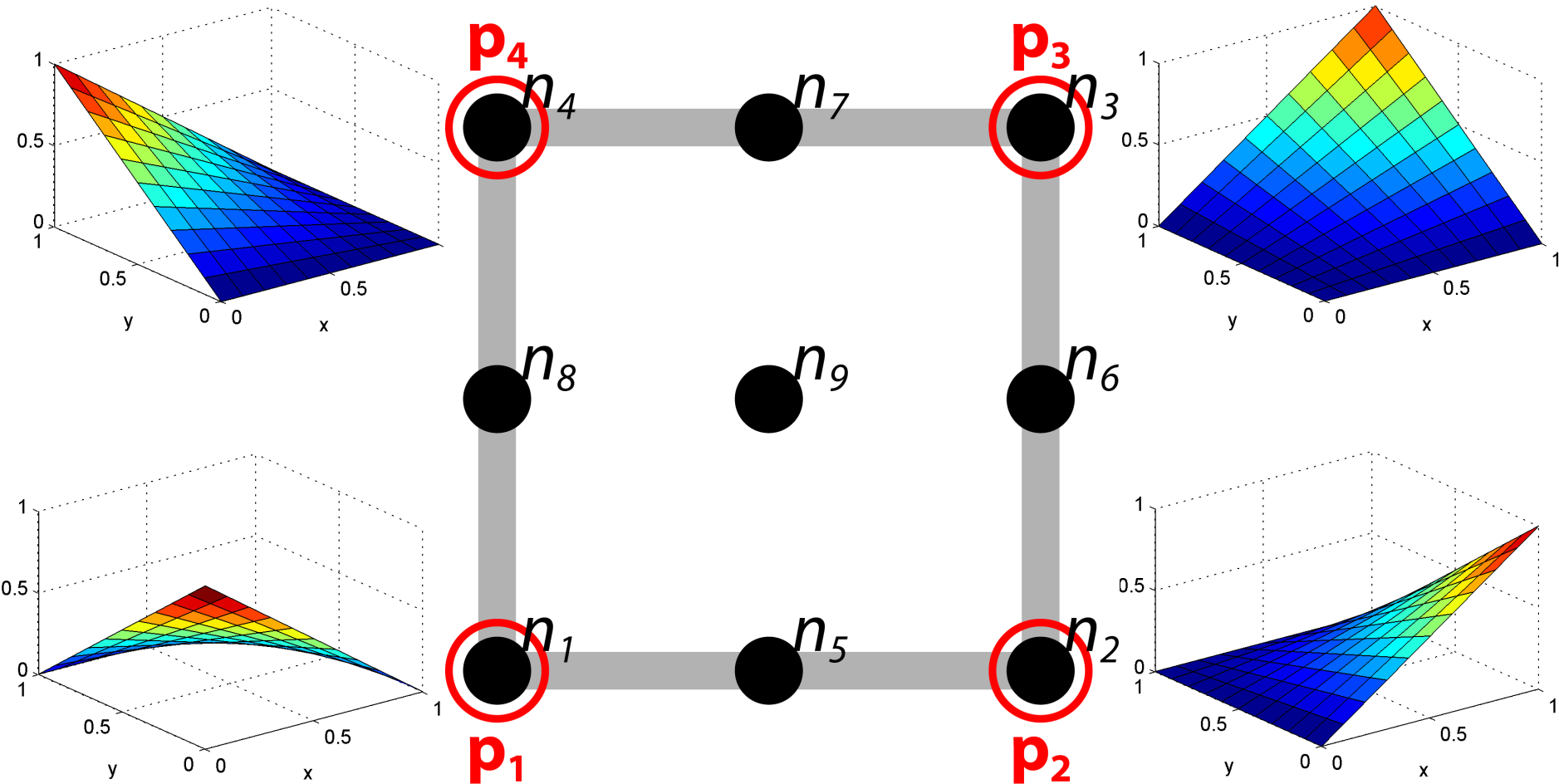
$$v_x(x, y) \approx N_1 v_{x,1} + N_2 v_{x,2} + \dots + N_9 v_{x,9}$$
$$= [N_1 \quad N_2 \quad \dots \quad N_9] \begin{Bmatrix} v_{x,1} \\ v_{x,2} \\ \vdots \\ v_{x,9} \end{Bmatrix} = \mathbf{N} \mathbf{v}_x$$

$$v_y(x, y) \approx \mathbf{N} \mathbf{v}_y$$

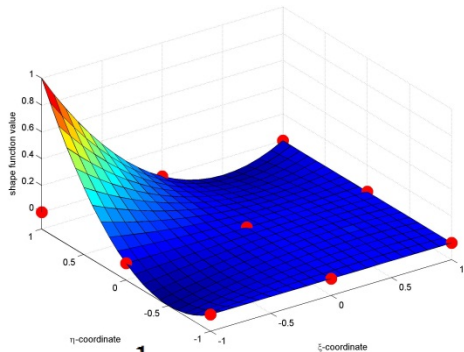
$$p(x, y) \approx N_1^p p_1 + N_2^p p_2 + N_3^p p_3 + N_4^p p_4$$

- 9 bi-quadratic shape functions for velocity
- 4 bi-linear shape functions for pressure (the same as for 2D elasticity)

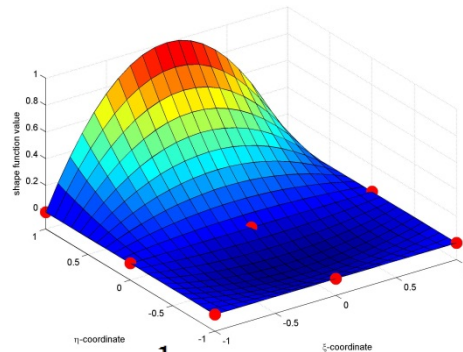
The Q2Q1-element: bi-linear pressure shape functions



The Q2Q1-element: bi-quadratic velocity shape functions

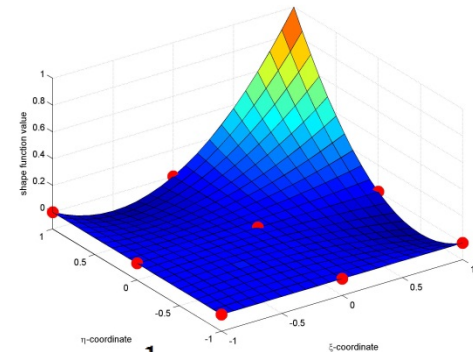


$$N_4 = \frac{1}{4}(\xi^2 - \xi)(\eta^2 + \eta)$$

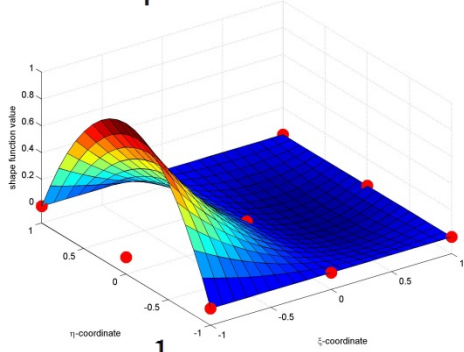


$$N_7 = -\frac{1}{2}(\xi^2 - 1)(\eta^2 + \eta)$$

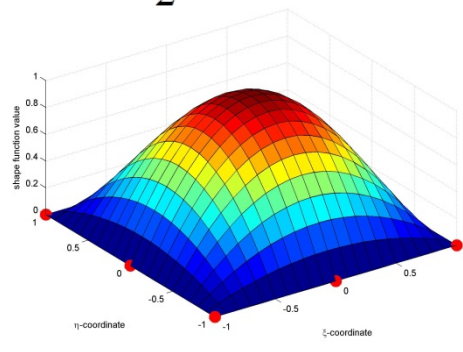
Element for
velocity



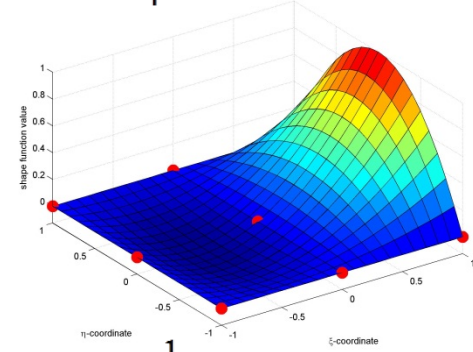
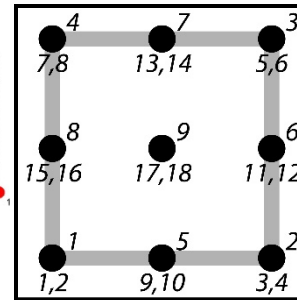
$$N_3 = \frac{1}{4}(\xi^2 + \xi)(\eta^2 + \eta)$$



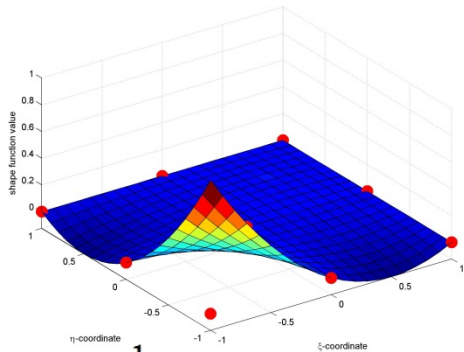
$$N_8 = -\frac{1}{2}(\xi^2 - \xi)(\eta^2 - 1)$$



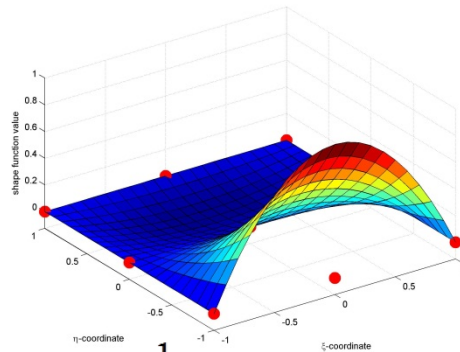
$$N_9 = (\xi^2 - 1)(\eta^2 - 1)$$



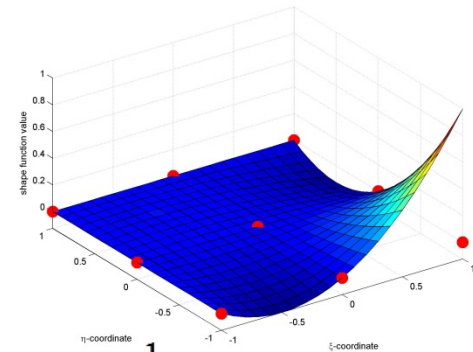
$$N_6 = -\frac{1}{2}(\xi^2 + \xi)(\eta^2 - 1)$$



$$N_1 = \frac{1}{4}(\xi^2 - \xi)(\eta^2 - \eta)$$



$$N_5 = -\frac{1}{2}(\xi^2 - 1)(\eta^2 - \eta)$$



$$N_2 = \frac{1}{4}(\xi^2 + \xi)(\eta^2 - \eta)$$

Force balance equations: The FE-approximation

④ Apply finite element approximation

$$\int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \left[\frac{4}{3} \eta \frac{\partial \mathbf{N}}{\partial x} \mathbf{v}_x - \frac{2}{3} \eta \frac{\partial \mathbf{N}}{\partial y} \mathbf{v}_y \right] dy dx + \int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \eta \left(\frac{\partial \mathbf{N}}{\partial y} \mathbf{v}_x + \frac{\partial \mathbf{N}}{\partial x} \mathbf{v}_y \right) dx dy - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \mathbf{N}_p \mathbf{p} dx dy = \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_x dx dy$$

$$\int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \left[-\frac{2}{3} \eta \frac{\partial \mathbf{N}}{\partial x} \mathbf{v}_x + \frac{4}{3} \eta \frac{\partial \mathbf{N}}{\partial y} \mathbf{v}_y \right] dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \eta \left(\frac{\partial \mathbf{N}}{\partial y} \mathbf{v}_x + \frac{\partial \mathbf{N}}{\partial x} \mathbf{v}_y \right) dy dx - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial y} \mathbf{N}_p \mathbf{p} dx dy = \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_y dx dy$$

- Take nodal values out of integration

$$\int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \frac{4}{3} \eta \frac{\partial \mathbf{N}}{\partial x} dy dx \mathbf{v}_x - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \frac{2}{3} \eta \frac{\partial \mathbf{N}}{\partial y} dy dx \mathbf{v}_y + \int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \eta \frac{\partial \mathbf{N}}{\partial y} dx dy \mathbf{v}_y$$

$$+ \int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \eta \frac{\partial \mathbf{N}}{\partial x} dx dy \mathbf{v}_x - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \mathbf{N}_p dx dy \mathbf{p} = \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_x dx dy$$

$$- \int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \frac{2}{3} \eta \frac{\partial \mathbf{N}}{\partial x} dx dy \mathbf{v}_x + \int_0^{\Delta y} \int_0^{\Delta x} \frac{\partial \mathbf{N}^T}{\partial y} \frac{4}{3} \eta \frac{\partial \mathbf{N}}{\partial y} dx dy \mathbf{v}_y + \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \eta \frac{\partial \mathbf{N}}{\partial y} dy dx \mathbf{v}_y$$

$$+ \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial x} \eta \frac{\partial \mathbf{N}}{\partial x} dy dx \mathbf{v}_x - \int_0^{\Delta x} \int_0^{\Delta y} \frac{\partial \mathbf{N}^T}{\partial y} \mathbf{N}_p dx dy \mathbf{p} = \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}^T F_y dx dy$$

Force balance equation: Some reorganization

- Writing everything as one equation in vector-matrix form

$$\int_0^{\Delta x} \int_0^{\Delta y} \left[\begin{array}{cc} \frac{\partial \mathbf{N}^T}{\partial x} \frac{4}{3} \eta \frac{\partial \mathbf{N}}{\partial x} + \frac{\partial \mathbf{N}^T}{\partial y} \eta \frac{\partial \mathbf{N}}{\partial y} & \frac{\partial \mathbf{N}^T}{\partial x} \left(-\frac{2}{3} \eta \right) \frac{\partial \mathbf{N}}{\partial y} + \frac{\partial \mathbf{N}^T}{\partial y} \eta \frac{\partial \mathbf{N}}{\partial x} \\ \frac{\partial \mathbf{N}^T}{\partial y} \left(-\frac{2}{3} \eta \right) \frac{\partial \mathbf{N}}{\partial x} + \frac{\partial \mathbf{N}^T}{\partial x} \eta \frac{\partial \mathbf{N}}{\partial y} & \frac{\partial \mathbf{N}^T}{\partial y} \frac{4}{3} \eta \frac{\partial \mathbf{N}}{\partial y} + \frac{\partial \mathbf{N}^T}{\partial x} \eta \frac{\partial \mathbf{N}}{\partial x} \end{array} \right] dx dy \begin{Bmatrix} v_{x,1} \\ \vdots \\ v_{x,9} \\ v_{y,1} \\ \vdots \\ v_{y,9} \end{Bmatrix} - \int_0^{\Delta x} \int_0^{\Delta y} \begin{Bmatrix} \frac{\partial \mathbf{N}^T}{\partial x} \\ \frac{\partial \mathbf{N}^T}{\partial y} \end{Bmatrix} \mathbf{N}_p dx dy \begin{Bmatrix} p_1 \\ \vdots \\ p_4 \end{Bmatrix} = \dots$$

- Reorganization

$$\int_0^{\Delta x} \int_0^{\Delta y} \begin{Bmatrix} \frac{\partial N_1}{\partial x} & 0 & \frac{\partial N_1}{\partial y} \\ 0 & \frac{\partial N_1}{\partial y} & \frac{\partial N_1}{\partial x} \\ \vdots & \vdots & \vdots \\ \frac{\partial N_9}{\partial x} & 0 & \frac{\partial N_9}{\partial y} \\ 0 & \frac{\partial N_9}{\partial y} & \frac{\partial N_9}{\partial x} \end{Bmatrix} \begin{Bmatrix} \frac{4}{3} \eta & -\frac{2}{3} \eta & 0 \\ -\frac{2}{3} \eta & \frac{4}{3} \eta & 0 \\ 0 & 0 & \eta \end{Bmatrix} \begin{Bmatrix} \frac{\partial N_1}{\partial x} & 0 & \dots & \frac{\partial N_9}{\partial x} & 0 \\ 0 & \frac{\partial N_1}{\partial y} & \dots & 0 & \frac{\partial N_9}{\partial y} \\ \frac{\partial N_1}{\partial y} & \frac{\partial N_1}{\partial x} & \dots & \frac{\partial N_9}{\partial y} & \frac{\partial N_9}{\partial x} \end{Bmatrix} dx dy \begin{Bmatrix} v_{x,1} \\ v_{y,1} \\ \vdots \\ v_{x,9} \\ v_{y,9} \end{Bmatrix} - \int_0^{\Delta x} \int_0^{\Delta y} \begin{Bmatrix} \frac{\partial N_1}{\partial x} \\ \frac{\partial N_1}{\partial y} \\ \vdots \\ \frac{\partial N_9}{\partial x} \\ \frac{\partial N_9}{\partial y} \end{Bmatrix} \mathbf{N}_p dx dy \begin{Bmatrix} p_1 \\ \vdots \\ p_4 \end{Bmatrix} = \dots$$

Force balance equation: The final equation

$$\int_0^{\Delta x} \int_0^{\Delta y} \underbrace{\begin{bmatrix} \partial_x N_1 & 0 & \partial_y N_1 \\ 0 & \partial_y N_1 & \partial_x N_1 \\ \vdots & \vdots & \vdots \\ \partial_x N_9 & 0 & \partial_y N_9 \\ 0 & \partial_y N_9 & \partial_x N_9 \end{bmatrix}}_{\mathbf{B}^T} \underbrace{\begin{bmatrix} \frac{4}{3}\eta & -\frac{2}{3}\eta & 0 \\ -\frac{2}{3}\eta & \frac{4}{3}\eta & 0 \\ 0 & 0 & \eta \end{bmatrix}}_{\mathbf{D}} \underbrace{\begin{bmatrix} \frac{\partial N_1}{\partial x} & 0 & \dots & \frac{\partial N_9}{\partial x} & 0 \\ 0 & \frac{\partial N_1}{\partial y} & \dots & 0 & \frac{\partial N_9}{\partial y} \\ \frac{\partial N_1}{\partial y} & \frac{\partial N_1}{\partial x} & \dots & \frac{\partial N_9}{\partial y} & \frac{\partial N_9}{\partial x} \end{bmatrix}}_{\mathbf{B}} dxdy \underbrace{\begin{Bmatrix} v_{x,1} \\ v_{y,1} \\ \vdots \\ v_{x,9} \\ v_{y,9} \end{Bmatrix}}_{\mathbf{v}}$$

K

$$- \int_0^{\Delta x} \int_0^{\Delta y} \underbrace{\begin{Bmatrix} \frac{\partial N_1}{\partial x} \\ \frac{\partial N_1}{\partial y} \\ \vdots \\ \frac{\partial N_9}{\partial x} \\ \frac{\partial N_9}{\partial y} \end{Bmatrix}}_{\mathbf{G}} \mathbf{N}_p dxdy \underbrace{\begin{Bmatrix} p_1 \\ \vdots \\ p_4 \end{Bmatrix}}_{\mathbf{p}} = \int_0^{\Delta x} \int_0^{\Delta y} \underbrace{\begin{Bmatrix} N_1 F_x \\ N_1 F_y \\ \vdots \\ N_9 F_x \\ N_9 F_y \end{Bmatrix}}_{\mathbf{F}} dxdy$$

$$\mathbf{Kv} - \mathbf{Gp} = \mathbf{F}$$

Apply numerical integration using 9 integration points

Incompressibility equation: Steps ① – ④

- Original equation:

$$\frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} = 0$$

- Test functions (use pressure functions):

$$\mathbf{N}_p = [N_1^p \quad \dots \quad N_4^p]$$

- ① Multiplication with test functions:

$$\mathbf{N}_p^T \frac{\partial v_x}{\partial x} + \mathbf{N}_p^T \frac{\partial v_y}{\partial y} = 0$$

- ② Spatial integration:

$$\int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}_p^T \frac{\partial v_x}{\partial x} dx dy + \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}_p^T \frac{\partial v_y}{\partial y} dx dy = 0$$

- ③ NO integration by parts!!!

- ④ Apply FE-approximation:

$$\int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}_p^T \frac{\partial \mathbf{N}}{\partial x} dx dy \mathbf{v}_x + \int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}_p^T \frac{\partial \mathbf{N}}{\partial y} dx dy \mathbf{v}_y = 0$$

Incompressibility equation: Some reorganization

- Writing everything as one equation in vector-matrix form

$$\int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}_p^T \begin{bmatrix} \frac{\partial \mathbf{N}}{\partial x} & \frac{\partial \mathbf{N}}{\partial y} \end{bmatrix} dx dy \begin{Bmatrix} v_{x,1} \\ \vdots \\ v_{x,9} \\ v_{y,1} \\ \vdots \\ v_{y,9} \end{Bmatrix} = 0$$

- Reorganization

$$\int_0^{\Delta x} \int_0^{\Delta y} \mathbf{N}_p^T \begin{bmatrix} \frac{\partial N_1}{\partial x} & \frac{\partial N_1}{\partial y} & \dots & \frac{\partial N_9}{\partial x} & \frac{\partial N_9}{\partial y} \end{bmatrix} dx dy \begin{Bmatrix} v_{x,1} \\ v_{y,1} \\ \vdots \\ v_{x,9} \\ v_{y,9} \end{Bmatrix} = 0$$

Apply numerical integration using 9 integration points

$$\mathbf{G}^T \mathbf{v} = \mathbf{0}$$

Final set of equations

- Force balance:

$$\mathbf{K}\mathbf{v} - \mathbf{G}\mathbf{p} = \mathbf{F}$$

- Incompressibility:

$$\mathbf{G}^T \mathbf{v} = \mathbf{0}$$

- Everything together:

$$\begin{bmatrix} \mathbf{K} & -\mathbf{G} \\ \mathbf{G}^T & \mathbf{0} \end{bmatrix} \begin{Bmatrix} \mathbf{v} \\ \mathbf{p} \end{Bmatrix} = \begin{Bmatrix} \mathbf{F} \\ \mathbf{0} \end{Bmatrix}$$

- Voilà:

$$\hat{\mathbf{K}}\mathbf{a} = \hat{\mathbf{F}}$$

What's new?

- New physical parameters (viscosity instead of elastic parameters)
- Numerical grid and indexing is different because of bi-quadratic shape functions (9 nodes per element).

So, you need to change

- `EL_N`
- `EL_DOF`

You need to newly introduce

- `EL_P`, which provides a relationship between element numbers and pressure-node numbers
- New set of shape functions (4 bi-linear for pressure stay the same, 9 bi-quadratic for velocity are new)
- New set of integration points (9 instead of 4)
- Calculate **G** inside the loop over integration points in a similar way as **K**. For this you need the shape function derivatives in a vector format (instead of matrix **B**).
- Put **K**, **G**, and **F** together to form the super-big **K** and super-big **F**.
- Your solution is now a velocity. So, introduce a time increment Δt and update your coordinates as **GCOORD** = **GCOORD** + Δt ***velocity**. Then you can write a time-loop around your code and run it for several time steps.

Problem to be solved

- Homogeneous medium
- Boundary conditions
 - Left and bottom: Free slip (i.e., no boundary-perpendicular velocity, boundary-parallel velocity left free)
 - Right and top: Free surface (both velocities left free)
- Calculate several time steps

